

Intro. to Fourier Series (10.2)

(F1)

Have already seen spaces of solutions ^(functs.) to certain ODE's form a finite dimal vect. space, lin. indep means any funct. in space can be written as lin. combs., i.e. they form a "basis" for the space.

Build upon this similarly to \mathbb{R}^n :
 in \mathbb{R}^3 : $\vec{v} = (x, y, z) = x\hat{i} + y\hat{j} + z\hat{k}$, $\{\hat{i}, \hat{j}, \hat{k}\}$ an orthonormal basis: $\hat{i} \cdot \hat{i} = 1$, $\hat{i} \cdot \hat{j} = \hat{i} \cdot \hat{k} = 0$, etc. and $\vec{v} = (\vec{v} \cdot \hat{i})\hat{i} + (\vec{v} \cdot \hat{j})\hat{j} + (\vec{v} \cdot \hat{k})\hat{k}$, $\forall \vec{v} \in \mathbb{R}^3$
 in \mathbb{R}^n : $\hat{e}_i = (0, \dots, 0, 1, 0, \dots)$ similar, inner (or dot) prod. $\vec{v} \cdot \hat{e}_i \rightarrow i^{\text{th}}$ component, i.e. $\vec{v} = \sum (\vec{v} \cdot \hat{e}_i) \hat{e}_i$

For fund. spaces, the basis vects. are functs. ϕ_i and the inner prod. of 2 funcs. $f + g$ is given by $\int_a^b f(x)g(x)dx$

$\{\phi_i\}$ as basis: orthog. cond. makes lin. indep. $\int_a^b \phi_i \phi_j = \delta_{ij} = \begin{cases} 0 & i \neq j \\ 1 & i = j \end{cases}$ (normal as well)

and any fund. f in the space can be written as $f = \sum_i \underbrace{\langle f, \phi_i \rangle}_{\text{Fourier coeff}} \phi_i$, the Fourier series of f

The classical one uses $\{e^{-inx}\}_{i=0}^{\infty} = \{\phi_i\}$, eqv. to $\{\cos \frac{m\pi x}{L}\} \cup \{\sin \frac{m\pi x}{L}\}$. [wavelets, etc.]

These are particularly useful when dealing w/ per. funcs. since they all are periodic

F2

Tripe. fnc. $\sin x + \cos x$ have
fund. per. 2π (see blc.), so $\frac{m\pi x}{L} = 2\pi \Rightarrow$

$$T = \frac{2L}{m} \text{ is fund. per. of } \sin + \cos \frac{m\pi x}{L}$$

(any integral mult. of T is also a period)

Write series as $f(x) = \frac{a_0}{2} + \sum_{m=1}^{\infty} \left(a_m \cos \frac{m\pi x}{L} + b_m \sin \frac{m\pi x}{L} \right)$,
using inner prod. $\frac{1}{L} \int_{-L}^L f(x)g(x)dx = \langle f, g \rangle$, see

$\langle \phi_m, \psi_n \rangle = \delta_{mn}$ by doing integrations
(recalling tripe. iden. + intep. by parts in blc)

Given a convergent Fourier series $f(x)$, the book
derives why the coeffs $a_m + b_m$ are what I claim:

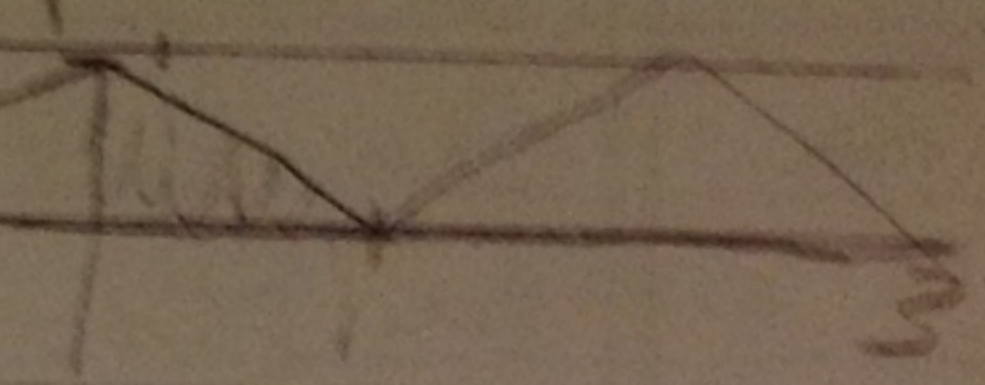
$$a_0 = \frac{1}{L} \int_{-L}^L f(x) dx, \quad a_n = \frac{1}{L} \int_{-L}^L f(x) \cos \frac{n\pi x}{L} dx, \quad n \geq 1$$

$$b_n = \frac{1}{L} \int_{-L}^L f(x) \sin \frac{n\pi x}{L} dx, \quad n \geq 1$$

Ex: p. 593 #16 assuming f converges:
 $f(x) = \begin{cases} x+1, & -1 \leq x < 0 \\ 1-x, & 0 \leq x < 1 \end{cases} \quad f(x+2) = f(x)$

So, fund. per. $T = 2L = 2, \quad L = 1$

$$a_0 = \int_{-1}^1 f(x) dx = 2 \int_0^1 = 2 \times \text{area} = 2 \left(\frac{1}{2} \right) = 1$$



(F3)

$$\begin{aligned}
 a_n &= \int_{-1}^1 f(x) \cos n\pi x \, dx = \int_{-1}^0 (1+x) \cos n\pi x \, dx + \int_0^1 (1-x) \cos n\pi x \, dx \\
 &= \int_{-1}^1 \cos n\pi x \, dx + 2 \int_0^1 (-x) \cos n\pi x \, dx \quad (\text{because of odd/even}) \\
 &\quad \text{by parts: } u=x, du=dx, dv=\cos n\pi x, v=\frac{1}{n\pi} \sin n\pi x \\
 &= \frac{1}{n\pi} \sin n\pi x \Big|_{-1}^1 - 2 \left(\frac{1}{n\pi} x \sin n\pi x + \frac{1}{(n\pi)^2} \cos n\pi x \right) \Big|_0^1 \\
 &= 0 + 2 \left(0 - 0 - \frac{1}{(n\pi)^2} [\cos n\pi - 1] \right), \quad n \geq 1 \\
 &= \frac{2}{(n\pi)^2} \begin{cases} 0 & n \text{ even} \\ 2 & n \text{ odd} \end{cases} = \begin{cases} 0, & n \text{ even} \\ \left(\frac{2}{n\pi}\right)^2, & n \text{ odd} \end{cases}
 \end{aligned}$$

b_n 's calc. similarly $\rightarrow b_n = 0, \forall n \geq 1$, so

$$\begin{aligned}
 f(x) &= \frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \cos n\pi x + b_n \sin n\pi x) \\
 &= \frac{1}{2} + \sum_{\substack{n \text{ odd} \\ n \geq 1}} \left(\frac{2}{n\pi}\right)^2 \cos n\pi x = \frac{1}{2} + \frac{4}{\pi^2} \sum_{n=1}^{\infty} \frac{\cos(2n-1)\pi x}{(2n-1)^2}
 \end{aligned}$$

Sec. 10.3 tells us when we get convergence:

Theorem (10.3.1) Suppose f & f' are piecewise cont. on $-L \leq x < L$ and f is periodic w/ per. $2L$ outside the int. $[-L, L)$, then the Fourier series for f converges to f at all pts. where f is cont., at the pts. of discont. $x=a$, it converges to $\frac{1}{2} \left(\lim_{x \rightarrow a^+} f(x) + \lim_{x \rightarrow a^-} f(x) \right) =$ avg. of left & rt. side limits at a

Recall f is piecewise cts. on $[a, b]$ if \exists part. of $[a, b]$ $a = x_0 < x_1 < x_2 < \dots < x_n = b$ s.t. f is cont. on ea. (x_{i-1}, x_i) and $\lim_{x \rightarrow x_i^+} f$ and $\lim_{x \rightarrow x_i^-} f$ exist & are finite

Note: for a = pt. of cty, $\lim_{x \rightarrow a^+} f = \lim_{x \rightarrow a^-} f = \lim_{x \rightarrow a} f$

So, Fourier series converges everywhere to $\frac{1}{2} (\lim_{x \rightarrow a^+} f(x) + \lim_{x \rightarrow a^-} f(x))$ [= f(a) if f is c.t.a.]

10.4 Recalls Even + Odd Funct.

odd $f(-x) = -f(x)$, even $f(-x) = f(x)$

$\int_{-L}^L \text{odd} = 0$
 $\int_{-L}^L \sin x$

$\int_{-L}^L \text{even} = 2 \int_0^L$
 $\int_{-L}^L \cos x$

prod of odd + even = odd

prod of 2 odd = even

sum of odds = odd

prod of 2 even = even

sum of evens = even

#13 p.608: prove any funct f can be written as sum of an even func. g and odd h:

Hint: what can you say about $f(x) + f(-x)$?

If $f(x) = g(x) + h(x)$, $f(-x) = g(x) - h(x)$
so $g(x) = \frac{1}{2} [f(x) + f(-x)] \Rightarrow h(x) = \frac{1}{2} [f(x) - f(-x)]$

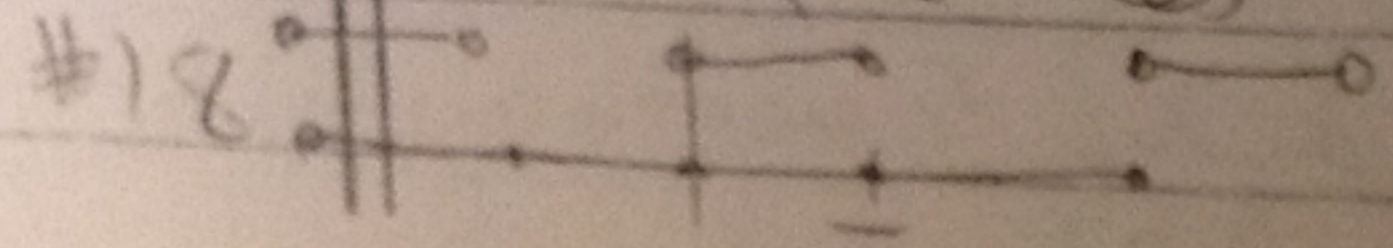
Given a func. $f(x)$ on $[0, L]$ can extend it to an odd or even func. w/ per. $2L$ by letting

even: $g(x) = \begin{cases} f(x) & 0 \leq x \leq L \\ f(-x) & -L < x < 0 \end{cases}$; odd: $h(x) = \begin{cases} f(x) & 0 < x < L \\ 0 & x = 0, L \\ -f(x) & -L < x < 0 \end{cases}$

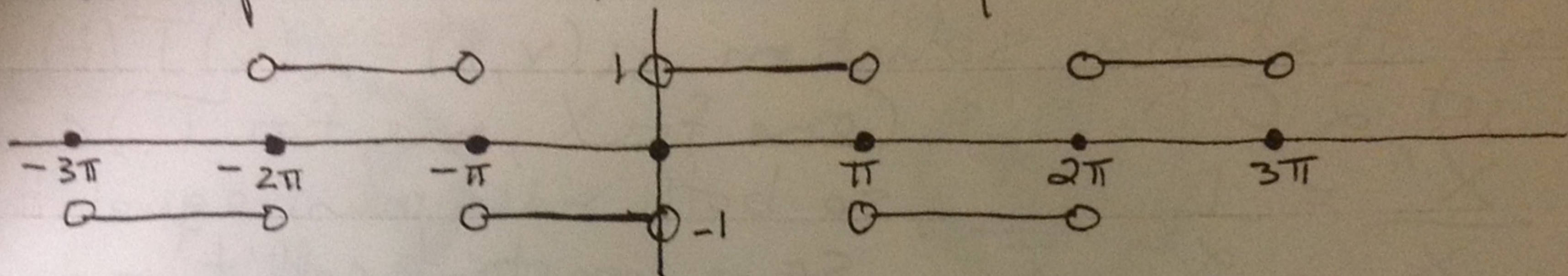
rep. by cosine series

rep. by sine series

Ex. #17 p.608 #17-18 $f(x) = 1$ on $[0, \pi]$ cosine series $f(x) = 1$



#18 on bottom of previous page (F4) was (F5) cut-off upon photocopying. Here it is again:
 The graph of the limiting function $F(x)$, which is the odd periodic extension of $f(x)=1$ on $(0, \pi)$ is



So the sine series for $f(x)=1$ on $(0, \pi)$ is given by $b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} F(x) \sin n x dx$ (general integrand)

$$= \frac{-2}{n\pi} \cos n x \Big|_0^{\pi} = \frac{2}{n\pi} (1 - \cos n\pi) = \begin{cases} 0, & n \text{ even} \\ \frac{4}{n\pi}, & n \text{ odd} \end{cases}$$

$\rightarrow \sum_{n=1}^{\infty} b_n \sin n x = \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{\sin(2n-1)x}{(2n-1)}$

Application to Simple P.D.E. (Heat Conduction 10.5)

Suppose you have a heat-conducting solid cylindrical rod whose length L is much greater than its diameter (eg. a wire), then the temperature in the rod is essentially a function of two independent variables: the distance along the rod x (ie. the spatial dimension along the axis of the cylinder) and the time t , call this function $u(x, t)$ and using the usual notation for partial derivatives, it turns out that

$$\alpha^2 u_{xx} = u_t \quad \text{for } 0 < x < L, t > 0$$

$\alpha^2 = \text{constant (called thermal diffusivity)}$

Assuming the initial condition $u(x,0) = f(x)$ (F6) on $0 \leq x \leq L$, and boundary conditions $u(0,t) = u(L,t) = 0$ for $t > 0$, along with a guess that the solution $u(x,t) = X(x)T(t)$, we get 2 O.D.E.'s (one for X , one for T) and $\frac{X''}{X} = \frac{1}{\alpha^2} \frac{T'}{T}$, so both sides must equal the same constant, call it $-\lambda$.

Solving these subject to the above boundary and initial conditions gives the solution

$$u(x,t) = \sum_{n=1}^{\infty} c_n e^{-n^2 \pi^2 \alpha^2 t / L^2} \sin\left(\frac{n\pi x}{L}\right)$$

where $c_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx$,

i.e. the c_n 's are the Fourier coefficients of the sine series for $f(x)$!

Ex. Solve $\alpha^2 u_{xx} = u_t$ for $0 \leq x \leq \pi$, $t > 0$

subject to $u(x,0) = 1$ and $u(0,t) = u(\pi,t) = 0$:

Using solution to #18 p. 608 (see page F4 + F5):

$$u(x,t) = \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{2n-1} e^{-(2n-1)^2 \alpha^2 t} \sin(2n-1)x$$